

NSF Workshop on the
Emerging Applications and Future Directions of the BEM

Fast Boundary Element Methods: A Mathematical View

Olaf Steinbach

Institute of Computational Mathematics, TU Graz, Austria

<http://www.numerik.math.tu-graz.ac.at>

NSF Workshop on the
Emerging Applications and Future Directions of the BEM

Fast Boundary Element Methods: A Mathematical View

Olaf Steinbach

Institute of Computational Mathematics, TU Graz, Austria

<http://www.numerik.math.tu-graz.ac.at>

This lecture is based on joint work with

- ▶ W. L. Wendland (Stuttgart), U. Langer (Linz), S. Rjasanow (Saarbrücken)
- ▶ G. Of; S. Engleder, G. Unger, P. Urthaler, M. Windisch, ...

Boundary element methods

- ▶ Direct vs. indirect formulation
- ▶ Weakly singular vs. hypersingular boundary integral equation
- ▶ 1st kind vs. 2nd kind boundary integral equation
- ▶ Collocation vs. Galerkin discretization
- ▶ pw constant vs. pw linear basis functions (hp BEM)
- ▶ Interpolation vs. projection of given boundary data
- ▶ Adaptive vs. uniform refinement
- ▶ Direct vs. preconditioned iterative solution (construction of preconditioners)
- ▶ Acceleration
(Panel Clustering, Fast Multipole, Adaptive Cross Approximation, ...)
- ▶ parallelization and domain decomposition methods
- ▶ ...

Boundary element methods

- ▶ Direct vs. indirect formulation
- ▶ Weakly singular vs. hypersingular boundary integral equation
- ▶ 1st kind vs. 2nd kind boundary integral equation
- ▶ Collocation vs. Galerkin discretization
- ▶ pw constant vs. pw linear basis functions (hp BEM)
- ▶ Interpolation vs. projection of given boundary data
- ▶ Adaptive vs. uniform refinement
- ▶ Direct vs. preconditioned iterative solution (construction of preconditioners)
- ▶ Acceleration
(Panel Clustering, Fast Multipole, Adaptive Cross Approximation, ...)
- ▶ parallelization and domain decomposition methods
- ▶ ...

While all of these topics are related to each other, we aim to end up with a most efficient, stable and accurate procedure to solve today's challenging problems from particular applications.

Model problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega \subset \mathbb{R}^3, \quad u(x) = g(x) \quad \text{for } x \in \Gamma = \partial\Omega$$

Model problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega \subset \mathbb{R}^3, \quad u(x) = g(x) \quad \text{for } x \in \Gamma = \partial\Omega$$

Indirect approach for $x \in \Omega$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} w(y) ds_y, \quad u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} v(y) ds_y$$

Model problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega \subset \mathbb{R}^3, \quad u(x) = g(x) \quad \text{for } x \in \Gamma = \partial\Omega$$

Indirect approach for $x \in \Omega$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} w(y) ds_y, \quad u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} v(y) ds_y$$

Boundary integral equations for $x \in \Gamma$

$$\frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} w(y) ds_y = g(x), \quad \frac{1}{2} v(x) + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} v(y) ds_y = g(x)$$

Model problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega \subset \mathbb{R}^3, \quad u(x) = g(x) \quad \text{for } x \in \Gamma = \partial\Omega$$

Indirect approach for $x \in \Omega$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} w(y) ds_y, \quad u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} v(y) ds_y$$

Boundary integral equations for $x \in \Gamma$

$$\frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} w(y) ds_y = g(x), \quad \frac{1}{2} v(x) + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} v(y) ds_y = g(x)$$

Representation formula for $x \in \Omega$ (direct approach)

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} \frac{\partial}{\partial n_y} u(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

Model problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega \subset \mathbb{R}^3, \quad u(x) = g(x) \quad \text{for } x \in \Gamma = \partial\Omega$$

Indirect approach for $x \in \Omega$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} w(y) ds_y, \quad u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} v(y) ds_y$$

Boundary integral equations for $x \in \Gamma$

$$\frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} w(y) ds_y = g(x), \quad \frac{1}{2} v(x) + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} v(y) ds_y = g(x)$$

Representation formula for $x \in \Omega$ (direct approach)

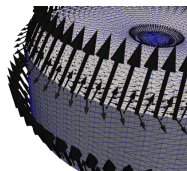
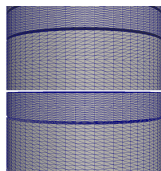
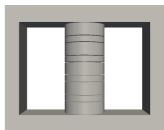
$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} \frac{\partial}{\partial n_y} u(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

Boundary integral equation for $x \in \Gamma$

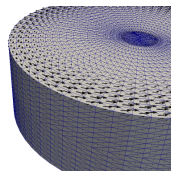
$$\frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y = \frac{1}{2} g(x) - \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} g(y) ds_y$$

Example: Controllable reactor [Z. Andjelic, G. Of, OS, P. Urthaler 2010]

- ▶ results in potential equation with piecewise constant material parameters
- ▶ indirect single layer potential approach
2nd kind boundary integral equation, simple, easy to solve
- ▶ direct domain decomposition approach (Steklov–Poincaré operator)
more advanced formulation
- ▶ on the continuous level both formulations are equivalent



indirect



direct

- ▶ difference in regularity of solutions → different approximation properties
- ▶ 122880 boundary elements → Adaptive Cross Approximation

Representation formula for $x \in \Omega$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

Boundary integral equation for $x \in \Gamma$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y + \frac{1}{2} u(x) + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

Representation formula for $x \in \Omega$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

Boundary integral equation for $x \in \Gamma$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y + \frac{1}{2} u(x) + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

Hypersingular boundary integral equation for $x \in \Gamma$

$$t(x) = \frac{1}{2} t(x) - \frac{1}{4\pi} \int_{\Gamma} \frac{(n_x, x-y)}{|x-y|^3} t(y) ds_y + \frac{1}{4\pi} \frac{\partial}{\partial n_x} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

Representation formula for $x \in \Omega$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

Boundary integral equation for $x \in \Gamma$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y + \frac{1}{2} u(x) + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

Hypersingular boundary integral equation for $x \in \Gamma$

$$t(x) = \frac{1}{2} t(x) - \frac{1}{4\pi} \int_{\Gamma} \frac{(n_x, x-y)}{|x-y|^3} t(y) ds_y + \frac{1}{4\pi} \frac{\partial}{\partial n_x} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} u(y) ds_y$$

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

Consequences

$$VK' = KV, \quad VD = \frac{1}{4}I - K^2$$

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

Consequences

$$VK' = KV, \quad VD = \frac{1}{4}I - K^2$$

- ▶ single layer potential symmetrizes the double layer potential [Plemelj 1911]
- ▶ single layer potential is a preconditioner for the hypersingular boundary integral operator, and vice versa [OS 1996; OS, Wendland 1998]

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

Consequences

$$VK' = KV, \quad VD = \frac{1}{4}I - K^2$$

- ▶ single layer potential symmetrizes the double layer potential [Plemelj 1911]
- ▶ single layer potential is a preconditioner for the hypersingular boundary integral operator, and vice versa [OS 1996; OS, Wendland 1998]

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$t = Su$$

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

Consequences

$$VK' = KV, \quad VD = \frac{1}{4}I - K^2$$

- ▶ single layer potential symmetrizes the double layer potential [Plemelj 1911]
- ▶ single layer potential is a preconditioner for the hypersingular boundary integral operator, and vice versa [OS 1996; OS, Wendland 1998]

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$t = Su = V^{-1}\left(\frac{1}{2}I + K\right)u$$

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

Consequences

$$VK' = KV, \quad VD = \frac{1}{4}I - K^2$$

- ▶ single layer potential symmetrizes the double layer potential [Plemelj 1911]
- ▶ single layer potential is a preconditioner for the hypersingular boundary integral operator, and vice versa [OS 1996; OS, Wendland 1998]

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$t = Su = V^{-1}\left(\frac{1}{2}I + K\right)u = \left[D + \left(\frac{1}{2}I + K'\right)V^{-1}\left(\frac{1}{2}I + K\right)\right]u = \dots$$

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

Consequences

$$VK' = KV, \quad VD = \frac{1}{4}I - K^2$$

- ▶ single layer potential symmetrizes the double layer potential [Plemelj 1911]
- ▶ single layer potential is a preconditioner for the hypersingular boundary integral operator, and vice versa [OS 1996; OS, Wendland 1998]

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$t = Su = V^{-1}\left(\frac{1}{2}I + K\right)u = \left[D + \left(\frac{1}{2}I + K'\right)V^{-1}\left(\frac{1}{2}I + K\right)\right]u = \dots$$

2nd kind boundary integral equation [OS, Wendland 2001]

$$\left(\frac{1}{2}I - K\right)v = g$$

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

Consequences

$$VK' = KV, \quad VD = \frac{1}{4}I - K^2$$

- ▶ single layer potential symmetrizes the double layer potential [Plemelj 1911]
- ▶ single layer potential is a preconditioner for the hypersingular boundary integral operator, and vice versa [OS 1996; OS, Wendland 1998]

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$t = Su = V^{-1}\left(\frac{1}{2}I + K\right)u = \left[D + \left(\frac{1}{2}I + K'\right)V^{-1}\left(\frac{1}{2}I + K\right)\right]u = \dots$$

2nd kind boundary integral equation [OS, Wendland 2001]

$$\left(\frac{1}{2}I - K\right)v = g, \quad v = \sum_{\ell=0}^{\infty} \left(\frac{1}{2}I + K\right)^{\ell} g$$

System of boundary integral equations (Calderon projector)

$$\begin{pmatrix} u \\ t \end{pmatrix} = \begin{pmatrix} \frac{1}{2}I - K & V \\ D & \frac{1}{2}I + K' \end{pmatrix} \begin{pmatrix} u \\ t \end{pmatrix}$$

Consequences

$$VK' = KV, \quad VD = \frac{1}{4}I - K^2$$

- ▶ single layer potential symmetrizes the double layer potential [Plemelj 1911]
- ▶ single layer potential is a preconditioner for the hypersingular boundary integral operator, and vice versa [OS 1996; OS, Wendland 1998]

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$t = Su = V^{-1}\left(\frac{1}{2}I + K\right)u = \left[D + \left(\frac{1}{2}I + K'\right)V^{-1}\left(\frac{1}{2}I + K\right)\right]u = \dots$$

2nd kind boundary integral equation [OS, Wendland 2001]

$$\left(\frac{1}{2}I - K\right)v = g, \quad v = \sum_{\ell=0}^{\infty} \left(\frac{1}{2}I + K\right)^{\ell} g, \quad \left\| \left(\frac{1}{2}I + K\right)v \right\|_{V^{-1}} \leq c_K \|v\|_{V^{-1}}, \quad c_K < 1$$

Boundary value problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Representation formula for $x \in \Omega$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} g(y) ds_y$$

Boundary value problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Representation formula for $x \in \Omega$

$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} g(y) ds_y$$

For an arbitrary boundary element solution t_h we define

$$\tilde{u}(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t_h(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} g(y) ds_y$$

Boundary value problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Representation formula for $x \in \Omega$

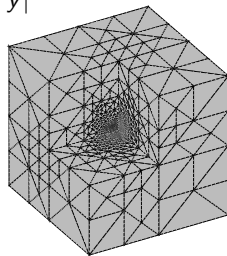
$$u(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} g(y) ds_y$$

For an arbitrary boundary element solution t_h we define

$$\tilde{u}(x) = \frac{1}{4\pi} \int_{\Gamma} \frac{1}{|x-y|} t_h(y) ds_y + \frac{1}{4\pi} \int_{\Gamma} \frac{(n_y, y-x)}{|x-y|^3} g(y) ds_y$$

Error equation for $x \in \Gamma$ [H. Schulz, OS 2000]

$$\left(\frac{1}{2}I - K'\right)[t - t_h](x) = \tilde{t}(x) - t_h(x)$$



Boundary value problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Boundary integral equation for $x \in \Gamma$

$$(Vt)(x) = \frac{1}{2}g(x) + (Kg)(x)$$

Approximation

$$t_h(x) = \sum_{k=1}^N t_k \psi_k(x), \quad g_h(x) = \sum_{i=1}^M g_i \varphi_i(x)$$

Boundary value problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Boundary integral equation for $x \in \Gamma$

$$(Vt)(x) = \frac{1}{2}g(x) + (Kg)(x)$$

Approximation

$$t_h(x) = \sum_{k=1}^N t_k \psi_k(x), \quad g_h(x) = \sum_{i=1}^M g_i \varphi_i(x)$$

Galerkin discretization

$$\sum_{k=1}^N t_k \int_{\Gamma} (V\psi_k)(x) \psi_{\ell}(x) ds_x = \sum_{i=1}^M g_i \int_{\Gamma} \left(\frac{1}{2}I + K\right) \varphi_i(x) \psi_{\ell}(x) ds_x \quad \text{for } \ell = 1, \dots, N$$

Boundary value problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Boundary integral equation for $x \in \Gamma$

$$(Vt)(x) = \frac{1}{2}g(x) + (Kg)(x)$$

Approximation

$$t_h(x) = \sum_{k=1}^N t_k \psi_k(x), \quad g_h(x) = \sum_{i=1}^M g_i \varphi_i(x)$$

Galerkin discretization

$$\sum_{k=1}^N t_k \int_{\Gamma} (V\psi_k)(x) \psi_{\ell}(x) ds_x = \sum_{i=1}^M g_i \int_{\Gamma} \left(\frac{1}{2}I + K\right) \varphi_i(x) \psi_{\ell}(x) ds_x \quad \text{for } \ell = 1, \dots, N$$

Algebraic system of linear equations

$$V_h \underline{t} = \left(\frac{1}{2}M_h + K_h\right) \underline{g}, \quad \underline{t} = V_h^{-1} \left(\frac{1}{2}M_h + K_h\right) \underline{g}$$

Discretization of Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg_h)(x) \varphi_j(x) ds_x = \int_{\Gamma} t_h(x) \varphi_j(x) ds_x = \sum_{k=1}^N t_k \int_{\Gamma} \psi_k(x) \varphi_j(x) ds_x$$

Discretization of Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg_h)(x) \varphi_j(x) ds_x = \int_{\Gamma} t_h(x) \varphi_j(x) ds_x = \sum_{k=1}^N t_k \int_{\Gamma} \psi_k(x) \varphi_j(x) ds_x$$

Boundary element approximation

$$S_h \underline{g} = M_h^{\top} \underline{t} = M_h^{\top} V_h^{-1} \left(\frac{1}{2} M_h + K_h \right) \underline{g}, \quad S_h = M_h^{\top} V_h^{-1} \left(\frac{1}{2} M_h + K_h \right)$$

Discretization of Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg_h)(x) \varphi_j(x) ds_x = \int_{\Gamma} t_h(x) \varphi_j(x) ds_x = \sum_{k=1}^N t_k \int_{\Gamma} \psi_k(x) \varphi_j(x) ds_x$$

Boundary element approximation

$$S_h \underline{g} = M_h^{\top} \underline{t} = M_h^{\top} V_h^{-1} \left(\frac{1}{2} M_h + K_h \right) \underline{g}, \quad S_h = M_h^{\top} V_h^{-1} \left(\frac{1}{2} M_h + K_h \right)$$

Stability (2D, ψ_k pw constant, ϖ_j pw linear)

$$M_h^{\top} = \frac{1}{2} h \begin{pmatrix} 1 & \cdots & 1 \\ 1 & 1 & \vdots \\ & 1 & \ddots & \vdots \\ & & \ddots & 1 \\ & & & 1 & 1 \end{pmatrix}$$

Discretization of Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg_h)(x) \varphi_j(x) ds_x = \int_{\Gamma} t_h(x) \varphi_j(x) ds_x = \sum_{k=1}^N t_k \int_{\Gamma} \psi_k(x) \varphi_j(x) ds_x$$

Boundary element approximation

$$S_h \underline{g} = M_h^{\top} \underline{t} = M_h^{\top} V_h^{-1} \left(\frac{1}{2} M_h + K_h \right) \underline{g}, \quad S_h = M_h^{\top} V_h^{-1} \left(\frac{1}{2} M_h + K_h \right)$$

Stability (2D, ψ_k pw constant, ϖ_j pw linear)

$$M_h^{\top} = \frac{1}{2} h \begin{pmatrix} 1 & \cdots & 1 \\ 1 & 1 & \vdots \\ & 1 & \ddots & \vdots \\ & & \ddots & 1 \\ & & & 1 & 1 \end{pmatrix}, \quad \underline{w} = \begin{pmatrix} 1 \\ -1 \\ \vdots \\ 1 \\ -1 \end{pmatrix}, \quad M_h^{\top} \underline{w} = \underline{0}$$

Discretization of Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg_h)(x) \varphi_j(x) ds_x = \int_{\Gamma} t_h(x) \varphi_j(x) ds_x = \sum_{k=1}^N t_k \int_{\Gamma} \psi_k(x) \varphi_j(x) ds_x$$

Boundary element approximation

$$S_h \underline{g} = M_h^T \underline{t} = M_h^T V_h^{-1} \left(\frac{1}{2} M_h + K_h \right) \underline{g}, \quad S_h = M_h^T V_h^{-1} \left(\frac{1}{2} M_h + K_h \right)$$

Stability (2D, ψ_k pw constant, ϖ_j pw linear)

$$M_h^T = \frac{1}{2} h \begin{pmatrix} 1 & \cdots & 1 \\ 1 & 1 & \vdots \\ & 1 & \ddots & \vdots \\ & & \ddots & 1 \\ & & & 1 & 1 \end{pmatrix}, \quad \underline{w} = \begin{pmatrix} 1 \\ -1 \\ \vdots \\ 1 \\ -1 \end{pmatrix}, \quad M_h^T \underline{w} = \underline{0}$$

Symmetric boundary element approximation

$$S_h = D_h + \left(\frac{1}{2} M_h^T + K_h^T \right) V_h^{-1} \left(\frac{1}{2} M_h + K_h \right)$$

Boundary value problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Variational formulation to find u , $u(x) = g(x)$ for $x \in \Gamma$:

$$\int_{\Omega} \nabla u(x) \cdot \nabla v(x) dx = 0 \quad \text{for all } v, v(x) = 0, x \in \Gamma$$

Boundary value problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Variational formulation to find u , $u(x) = g(x)$ for $x \in \Gamma$:

$$\int_{\Omega} \nabla u(x) \cdot \nabla v(x) dx = 0 \quad \text{for all } v, v(x) = 0, x \in \Gamma$$

Finite element formulation for $\ell = 1, \dots, M_{\Omega}$

$$\sum_{k=1}^{M_{\Omega}} u_k \int_{\Omega} \nabla \varphi_k(x) \cdot \nabla \varphi_{\ell}(x) dx + \sum_{i=M_{\Omega}+1}^M g_i \int_{\Omega} \nabla \varphi_i(x) \cdot \nabla \varphi_{\ell}(x) dx = 0$$

Linear system of algebraic equations

$$K_{II} \underline{u}_I + K_{CI} \underline{g} = \underline{0}$$

Boundary value problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Variational formulation to find u , $u(x) = g(x)$ for $x \in \Gamma$:

$$\int_{\Omega} \nabla u(x) \cdot \nabla v(x) dx = 0 \quad \text{for all } v, v(x) = 0, x \in \Gamma$$

Finite element formulation for $\ell = 1, \dots, M_{\Omega}$

$$\sum_{k=1}^{M_{\Omega}} u_k \int_{\Omega} \nabla \varphi_k(x) \cdot \nabla \varphi_{\ell}(x) dx + \sum_{i=M_{\Omega}+1}^M g_i \int_{\Omega} \nabla \varphi_i(x) \cdot \nabla \varphi_{\ell}(x) dx = 0$$

Linear system of algebraic equations

$$K_{II} \underline{u}_I + K_{CI} \underline{g} = \underline{0}, \quad \underline{u}_I = -K_{II}^{-1} K_{CI} \underline{g}$$

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg)(x)v(x)ds_x = \int_{\Gamma} \frac{\partial}{\partial n_x} u(x)v(x)ds_x$$

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg)(x)v(x)ds_x = \int_{\Gamma} \frac{\partial}{\partial n_x} u(x)v(x)ds_x = \int_{\Omega} \nabla u(x) \cdot \nabla v(x)dx$$

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg)(x)v(x)ds_x = \int_{\Gamma} \frac{\partial}{\partial n_x} u(x)v(x)ds_x = \int_{\Omega} \nabla u(x) \cdot \nabla v(x)dx$$

Discretization for $j = M_{\Omega} + 1, \dots, M$

$$\int_{\Gamma} Sg_h \varphi_j ds_x = \sum_{k=1}^{M_{\Omega}} u_k \int_{\Omega} \nabla \varphi_k \cdot \nabla \varphi_j dx + \sum_{i=M_{\Omega}+1}^M g_i \int_{\Omega} \nabla \varphi_i \cdot \nabla \varphi_j dx$$

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg)(x)v(x)ds_x = \int_{\Gamma} \frac{\partial}{\partial n_x} u(x)v(x)ds_x = \int_{\Omega} \nabla u(x) \cdot \nabla v(x)dx$$

Discretization for $j = M_{\Omega} + 1, \dots, M$

$$\int_{\Gamma} Sg_h \varphi_j ds_x = \sum_{k=1}^{M_{\Omega}} u_k \int_{\Omega} \nabla \varphi_k \cdot \nabla \varphi_j dx + \sum_{i=M_{\Omega}+1}^M g_i \int_{\Omega} \nabla \varphi_i \cdot \nabla \varphi_j dx$$

Finite element approximation

$$S_h \underline{g} = K_{IC} \underline{u}_I + K_{CC} \underline{g}$$

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg)(x)v(x)ds_x = \int_{\Gamma} \frac{\partial}{\partial n_x} u(x)v(x)ds_x = \int_{\Omega} \nabla u(x) \cdot \nabla v(x)dx$$

Discretization for $j = M_{\Omega} + 1, \dots, M$

$$\int_{\Gamma} Sg_h \varphi_j ds_x = \sum_{k=1}^{M_{\Omega}} u_k \int_{\Omega} \nabla \varphi_k \cdot \nabla \varphi_j dx + \sum_{i=M_{\Omega}+1}^M g_i \int_{\Omega} \nabla \varphi_i \cdot \nabla \varphi_j dx$$

Finite element approximation

$$S_h \underline{g} = K_{IC} \underline{u}_I + K_{CC} \underline{g} = [K_C - K_{IC} K_{II}^{-1} K_{CI}] \underline{g}$$

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg)(x)v(x)ds_x = \int_{\Gamma} \frac{\partial}{\partial n_x} u(x)v(x)ds_x = \int_{\Omega} \nabla u(x) \cdot \nabla v(x)dx$$

Discretization for $j = M_{\Omega} + 1, \dots, M$

$$\int_{\Gamma} Sg_h \varphi_j ds_x = \sum_{k=1}^{M_{\Omega}} u_k \int_{\Omega} \nabla \varphi_k \cdot \nabla \varphi_j dx + \sum_{i=M_{\Omega}+1}^M g_i \int_{\Omega} \nabla \varphi_i \cdot \nabla \varphi_j dx$$

Finite element approximation

$$S_h \underline{g} = K_{IC} \underline{u}_I + K_{CC} \underline{g} = [K_C - K_{IC} K_{II}^{-1} K_{CI}] \underline{g}$$

Discrete Steklov–Poincaré operators

$$S_h^{\text{BEM}} = D_h + \left(\frac{1}{2} M_h^{\top} + K_h^{\top}\right) V_h^{-1} \left(\frac{1}{2} M_h + K_h\right), \quad S_h^{\text{FEM}} = K_C - K_{IC} K_{II}^{-1} K_{CI}$$

Dirichlet to Neumann map (Steklov–Poincaré operator)

$$\int_{\Gamma} (Sg)(x)v(x)ds_x = \int_{\Gamma} \frac{\partial}{\partial n_x} u(x)v(x)ds_x = \int_{\Omega} \nabla u(x) \cdot \nabla v(x)dx$$

Discretization for $j = M_{\Omega} + 1, \dots, M$

$$\int_{\Gamma} Sg_h \varphi_j ds_x = \sum_{k=1}^{M_{\Omega}} u_k \int_{\Omega} \nabla \varphi_k \cdot \nabla \varphi_j dx + \sum_{i=M_{\Omega}+1}^M g_i \int_{\Omega} \nabla \varphi_i \cdot \nabla \varphi_j dx$$

Finite element approximation

$$S_h \underline{g} = K_{IC} \underline{u}_I + K_{CC} \underline{g} = [K_C - K_{IC} K_{II}^{-1} K_{CI}] \underline{g}$$

Discrete Steklov–Poincaré operators

$$S_h^{\text{BEM}} = D_h + \left(\frac{1}{2} M_h^{\top} + K_h^{\top}\right) V_h^{-1} \left(\frac{1}{2} M_h + K_h\right), \quad S_h^{\text{FEM}} = K_C - K_{IC} K_{II}^{-1} K_{CI}$$

- ▶ different FE/BE approximations of the Steklov–Poincaré operator
- ▶ coupling of FEM/BEM via domain decomposition methods

Model problem in 1D

$$-u''(x) = f(x) \quad \text{for } x \in (0, 1), \quad u(0) = u(1) = 0$$

Model problem in 1D

$$-u''(x) = f(x) \quad \text{for } x \in (0, 1), \quad u(0) = u(1) = 0$$

Finite difference approximation

$$h = \frac{1}{n}, \quad x_k = kh, \quad -u''(x_k) \approx \frac{-u_{k-1} + 2u_k - u_{k+1}}{h^2} \quad \text{for } k = 1, \dots, n-1$$

$$n = 9$$

$$K_9 = 81 \begin{pmatrix} 2 & -1 & & & & & & & \\ -1 & 2 & -1 & & & & & & \\ & -1 & 2 & -1 & & & & & \\ & & -1 & 2 & -1 & & & & \\ & & & -1 & 2 & -1 & & & \\ & & & & -1 & 2 & -1 & & \\ & & & & & -1 & 2 & -1 & \\ & & & & & & -1 & 2 & -1 \\ & & & & & & & -1 & 2 \end{pmatrix}$$

$$n = 9$$

$$K_9 = 81 \begin{pmatrix} 2 & -1 & & & & & & & \\ -1 & 2 & -1 & & & & & & \\ & -1 & 2 & -1 & & & & & \\ & & -1 & 2 & -1 & & & & \\ & & & -1 & 2 & -1 & & & \\ & & & & -1 & 2 & -1 & & \\ & & & & & -1 & 2 & -1 & \\ & & & & & & -1 & 2 & -1 \\ & & & & & & & -1 & 2 \end{pmatrix}$$

$$K_9^{-1} = \frac{1}{81} \cdot \frac{1}{9} \begin{pmatrix} 8 & 7 & 6 & 5 & 4 & 3 & 2 & 1 \\ 7 & 14 & 12 & 10 & 8 & 6 & 4 & 2 \\ 6 & 12 & 18 & 15 & 12 & 9 & 6 & 3 \\ 5 & 10 & 15 & 20 & 16 & 12 & 8 & 4 \\ 4 & 8 & 12 & 16 & 20 & 15 & 10 & 5 \\ 3 & 6 & 9 & 12 & 15 & 18 & 12 & 6 \\ 2 & 4 & 6 & 8 & 10 & 12 & 14 & 7 \\ 1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 \end{pmatrix}$$

$$n = 9$$

$$K_9 = 81 \begin{pmatrix} 2 & -1 & & & & & & & \\ -1 & 2 & -1 & & & & & & \\ & -1 & 2 & -1 & & & & & \\ & & -1 & 2 & -1 & & & & \\ & & & -1 & 2 & -1 & & & \\ & & & & -1 & 2 & -1 & & \\ & & & & & -1 & 2 & -1 & \\ & & & & & & -1 & 2 & -1 \\ & & & & & & & -1 & 2 \end{pmatrix}$$

$$K_9^{-1} = \frac{1}{81} \cdot \frac{1}{9} \begin{pmatrix} 8 & 7 & 6 & 5 & 4 & 3 & 2 & 1 \\ 7 & 14 & 12 & 10 & 8 & 6 & 4 & 2 \\ 6 & 12 & 18 & 15 & 12 & 9 & 6 & 3 \\ 5 & 10 & 15 & 20 & 16 & 12 & 8 & 4 \\ 4 & 8 & 12 & 16 & 20 & 15 & 10 & 5 \\ 3 & 6 & 9 & 12 & 15 & 18 & 12 & 6 \\ 2 & 4 & 6 & 8 & 10 & 12 & 14 & 7 \\ 1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 \end{pmatrix}$$

- ▶ inverse FDM/FEM stiffness matrix is dense,

$$n = 9$$

$$K_9 = 81 \begin{pmatrix} 2 & -1 & & & & & & & \\ -1 & 2 & -1 & & & & & & \\ & -1 & 2 & -1 & & & & & \\ & & -1 & 2 & -1 & & & & \\ & & & -1 & 2 & -1 & & & \\ & & & & -1 & 2 & -1 & & \\ & & & & & -1 & 2 & -1 & \\ & & & & & & -1 & 2 & -1 \\ & & & & & & & -1 & 2 \end{pmatrix}$$

$$K_9^{-1} = \frac{1}{81} \cdot \frac{1}{9} \begin{pmatrix} 8 & 7 & 6 & 5 & 4 & 3 & 2 & 1 \\ 7 & 14 & 12 & 10 & 8 & 6 & 4 & 2 \\ 6 & 12 & 18 & 15 & 12 & 9 & 6 & 3 \\ 5 & 10 & 15 & 20 & 16 & 12 & 8 & 4 \\ 4 & 8 & 12 & 16 & 20 & 15 & 10 & 5 \\ 3 & 6 & 9 & 12 & 15 & 18 & 12 & 6 \\ 2 & 4 & 6 & 8 & 10 & 12 & 14 & 7 \\ 1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 \end{pmatrix}$$

- ▶ inverse FDM/FEM stiffness matrix is dense, **but data sparse!**

$$K_9^{-1} = \frac{1}{729} \left(\begin{array}{cc} \begin{pmatrix} 8 & 7 \\ 7 & 14 \end{pmatrix} & \begin{pmatrix} 1 \\ 2 \end{pmatrix} \begin{pmatrix} 6 & 5 \end{pmatrix} \\ \begin{pmatrix} 6 \\ 5 \end{pmatrix} \begin{pmatrix} 1 & 2 \end{pmatrix} & \begin{pmatrix} 18 & 15 \\ 15 & 20 \end{pmatrix} \\ \begin{pmatrix} 4 \\ 3 \\ 2 \\ 1 \end{pmatrix} \begin{pmatrix} 1 & 2 & 3 & 4 \end{pmatrix} & \begin{pmatrix} 1 \\ 2 \\ 3 \\ 4 \end{pmatrix} \begin{pmatrix} 4 & 3 & 2 & 1 \end{pmatrix} \\ \begin{pmatrix} 20 & 15 \\ 15 & 18 \end{pmatrix} & \begin{pmatrix} 5 \\ 6 \end{pmatrix} \begin{pmatrix} 2 & 1 \end{pmatrix} \\ \begin{pmatrix} 2 \\ 1 \end{pmatrix} \begin{pmatrix} 5 & 6 \end{pmatrix} & \begin{pmatrix} 14 & 7 \\ 7 & 8 \end{pmatrix} \end{array} \right).$$

$$K_9^{-1} = \frac{1}{729} \left(\begin{array}{cc} \left(\begin{array}{cc} 8 & 7 \\ 7 & 14 \end{array} \right) & \left(\begin{array}{c} 1 \\ 2 \end{array} \right) \left(\begin{array}{cc} 6 & 5 \end{array} \right) \\ \left(\begin{array}{c} 6 \\ 5 \end{array} \right) \left(\begin{array}{cc} 1 & 2 \end{array} \right) & \left(\begin{array}{cc} 18 & 15 \\ 15 & 20 \end{array} \right) \\ \left(\begin{array}{c} 4 \\ 3 \\ 2 \\ 1 \end{array} \right) \left(\begin{array}{cccc} 1 & 2 & 3 & 4 \end{array} \right) & \left(\begin{array}{c} 1 \\ 2 \\ 3 \\ 4 \end{array} \right) \left(\begin{array}{cccc} 4 & 3 & 2 & 1 \end{array} \right) \\ \left(\begin{array}{c} 20 \\ 15 \\ 15 \\ 1 \end{array} \right) \left(\begin{array}{cc} 5 & 6 \end{array} \right) & \left(\begin{array}{c} 5 \\ 6 \end{array} \right) \left(\begin{array}{cc} 2 & 1 \end{array} \right) \\ \left(\begin{array}{c} 2 \\ 1 \end{array} \right) \left(\begin{array}{cc} 5 & 6 \end{array} \right) & \left(\begin{array}{cc} 14 & 7 \\ 7 & 8 \end{array} \right) \end{array} \right).$$

Storage requirement for off diagonal block

$$\left(\frac{n-1}{2}\right)^2 = \frac{1}{4}(n-1)^2 \quad \rightarrow \quad 2 \frac{n-1}{2} = n-1$$

- ▶ concept of hierarchical matrices [Hackbusch, ...]
 - hierarchical block decomposition of the matrix
 - low rank approximation of blocks

- ▶ concept of hierarchical matrices [Hackbusch, ...]
 - hierarchical block decomposition of the matrix
 - low rank approximation of blocks
- ▶ **But why does this works?**

- ▶ concept of hierarchical matrices [Hackbusch, ...]
 - hierarchical block decomposition of the matrix
 - low rank approximation of blocks
- ▶ **But why does this works?**

Boundary value problem

$$-u''(x) = f(x) \quad \text{for } x \in (0, 1), \quad u(0) = u(1) = 0$$

- ▶ concept of hierarchical matrices [Hackbusch, ...]
 - hierarchical block decomposition of the matrix
 - low rank approximation of blocks
- ▶ **But why does this works?**

Boundary value problem

$$-u''(x) = f(x) \quad \text{for } x \in (0, 1), \quad u(0) = u(1) = 0$$

Solution

$$u(x) = \int_0^1 G(x, y) f(y) dy, \quad G(x, y) = \begin{cases} (1-x)y & \text{for } 0 < y < x, \\ x(1-y) & \text{for } x < y < 1. \end{cases}$$

- ▶ concept of hierarchical matrices [Hackbusch, ...]
 - hierarchical block decomposition of the matrix
 - low rank approximation of blocks
- ▶ **But why does this works?**

Boundary value problem

$$-u''(x) = f(x) \quad \text{for } x \in (0, 1), \quad u(0) = u(1) = 0$$

Solution

$$u(x) = \int_0^1 G(x, y) f(y) dy, \quad G(x, y) = \begin{cases} (1-x)y & \text{for } 0 < y < x, \\ x(1-y) & \text{for } x < y < 1. \end{cases}$$

Approximation

$$u_k \approx u(x_k) = \int_0^1 G(x_k, y) f(y) dy \approx \sum_{\ell=1}^n \omega_{\ell} G(x_k, y_{\ell}) f(y_{\ell})$$

- ▶ concept of hierarchical matrices [Hackbusch, ...]
 - hierarchical block decomposition of the matrix
 - low rank approximation of blocks
- ▶ **But why does this works?**

Boundary value problem

$$-u''(x) = f(x) \quad \text{for } x \in (0, 1), \quad u(0) = u(1) = 0$$

Solution

$$u(x) = \int_0^1 G(x, y) f(y) dy, \quad G(x, y) = \begin{cases} (1-x)y & \text{for } 0 < y < x, \\ x(1-y) & \text{for } x < y < 1. \end{cases}$$

Approximation

$$u_k \approx u(x_k) = \int_0^1 G(x_k, y) f(y) dy \approx \sum_{\ell=1}^n \omega_\ell G(x_k, y_\ell) f(y_\ell), \quad \underline{u} = A_n \underline{f}$$

with

$$A_n[\ell, k] = \omega_\ell G(x_k, y_\ell) \omega_\ell$$

- ▶ concept of hierarchical matrices [Hackbusch, ...]
 - hierarchical block decomposition of the matrix
 - low rank approximation of blocks
- ▶ **But why does this works?**

Boundary value problem

$$-u''(x) = f(x) \quad \text{for } x \in (0, 1), \quad u(0) = u(1) = 0$$

Solution

$$u(x) = \int_0^1 G(x, y) f(y) dy, \quad G(x, y) = \begin{cases} (1-x)y & \text{for } 0 < y < x, \\ x(1-y) & \text{for } x < y < 1. \end{cases}$$

Approximation

$$u_k \approx u(x_k) = \int_0^1 G(x_k, y) f(y) dy \approx \sum_{\ell=1}^n \omega_\ell G(x_k, y_\ell) f(y_\ell), \quad \underline{u} = A_n \underline{f}$$

with

$$A_n[\ell, k] = \omega_\ell G(x_k, y_\ell) \omega_k = (1-x_k) y_\ell \omega_\ell \quad \text{for all } x_k < y_\ell$$

- ▶ concept of hierarchical matrices [Hackbusch, ...]
 - hierarchical block decomposition of the matrix
 - low rank approximation of blocks
- ▶ **But why does this works?**

Boundary value problem

$$-u''(x) = f(x) \quad \text{for } x \in (0, 1), \quad u(0) = u(1) = 0$$

Solution

$$u(x) = \int_0^1 G(x, y) f(y) dy, \quad G(x, y) = \begin{cases} (1-x)y & \text{for } 0 < y < x, \\ x(1-y) & \text{for } x < y < 1. \end{cases}$$

Approximation

$$u_k \approx u(x_k) = \int_0^1 G(x_k, y) f(y) dy \approx \sum_{\ell=1}^n \omega_\ell G(x_k, y_\ell) f(y_\ell), \quad \underline{u} = A_n \underline{f}$$

with

$$A_n[\ell, k] = \omega_\ell G(x_k, y_\ell) \omega_\ell = (1-x_k) y_\ell \omega_\ell \quad \text{for all } x_k < y_\ell$$

→ Rank 1 representation of matrix block

Low rank approximation of matrices

$$A \in \mathbb{R}^{m \times n}, \quad \mu = \text{rank } A \leq \min\{m, n\}$$

with

$$A^T A \in \mathbb{R}^{n \times n}, \quad \lambda_1(A^T A) \geq \lambda_2(A^T A) \geq \dots \geq \lambda_n(A^T A) \geq 0$$

Low rank approximation of matrices

$$A \in \mathbb{R}^{m \times n}, \quad \mu = \text{rank } A \leq \min\{m, n\}$$

with

$$A^T A \in \mathbb{R}^{n \times n}, \quad \lambda_1(A^T A) \geq \lambda_2(A^T A) \geq \dots \geq \lambda_n(A^T A) \geq 0$$

Singular values

$$\sigma_k := \sqrt{\lambda_k(A^T A)} \geq 0 \quad \text{for } k = 1, \dots, \mu$$

Low rank approximation of matrices

$$A \in \mathbb{R}^{m \times n}, \quad \mu = \text{rank } A \leq \min\{m, n\}$$

with

$$A^T A \in \mathbb{R}^{n \times n}, \quad \lambda_1(A^T A) \geq \lambda_2(A^T A) \geq \dots \geq \lambda_n(A^T A) \geq 0$$

Singular values

$$\sigma_k := \sqrt{\lambda_k(A^T A)} \geq 0 \quad \text{for } k = 1, \dots, \mu$$

Singular value decomposition

$$D := \text{diag}(\lambda_k(A)) = V^T A^T A V, \quad A^T A \underline{v}_k = \lambda_k \underline{v}_k$$

Low rank approximation of matrices

$$A \in \mathbb{R}^{m \times n}, \quad \mu = \text{rank } A \leq \min\{m, n\}$$

with

$$A^T A \in \mathbb{R}^{n \times n}, \quad \lambda_1(A^T A) \geq \lambda_2(A^T A) \geq \dots \geq \lambda_n(A^T A) \geq 0$$

Singular values

$$\sigma_k := \sqrt{\lambda_k(A^T A)} \geq 0 \quad \text{for } k = 1, \dots, \mu$$

Singular value decomposition

$$D := \text{diag}(\lambda_k(A)) = V^T A^T A V, \quad A^T A \underline{v}_k = \lambda_k \underline{v}_k$$

$$\Sigma = \text{diag}(\sigma_k(A)) \in \mathbb{R}^{m \times n}, \quad D = \Sigma^T \Sigma, \quad U = A V \Sigma^+$$

Low rank approximation of matrices

$$A \in \mathbb{R}^{m \times n}, \quad \mu = \text{rank } A \leq \min\{m, n\}$$

with

$$A^T A \in \mathbb{R}^{n \times n}, \quad \lambda_1(A^T A) \geq \lambda_2(A^T A) \geq \dots \geq \lambda_n(A^T A) \geq 0$$

Singular values

$$\sigma_k := \sqrt{\lambda_k(A^T A)} \geq 0 \quad \text{for } k = 1, \dots, \mu$$

Singular value decomposition

$$D := \text{diag}(\lambda_k(A)) = V^T A^T A V, \quad A^T A \underline{v}_k = \lambda_k \underline{v}_k$$

$$\Sigma = \text{diag}(\sigma_k(A)) \in \mathbb{R}^{m \times n}, \quad D = \Sigma^T \Sigma, \quad U = A V \Sigma^+$$

$$A = U \Sigma V^T = \sum_{k=1}^{\mu} \sigma_k(A) \underline{u}_k \underline{v}_k^T$$

Approximation

$$A_r = \sum_{k=1}^r \sigma_k(A) \underline{u}_k \underline{v}_k^\top, \quad \|A - A_r\|_2 = \sigma_{r+1}(A)$$

Approximation

$$A_r = \sum_{k=1}^r \sigma_k(A) \underline{u}_k \underline{v}_k^\top, \quad \|A - A_r\|_2 = \sigma_{r+1}(A)$$

Storage requirement

$$nm \quad \text{vs} \quad r(n+m)$$

Matrix by vector multiplication

$$A_r \underline{w} = \sum_{k=1}^r \sigma_k(A) \underline{u}_k \underline{v}_k^\top \underline{w} = \sum_{k=1}^r \sigma_k(A) (\underline{v}_k^\top \underline{w}) \underline{u}_k, \quad r(n+m)$$

- ▶ low rank matrix is singular
- ▶ block decomposition due to some admissibility condition
- ▶ low rank approximation of blocks

Model problem

$$-\Delta u(x) = 0 \quad \text{for } x \in \Omega, \quad u(x) = g(x) \quad \text{for } x \in \Gamma$$

Boundary integral equation for $x \in \Gamma$

$$\int_{\Gamma} U^*(x, y) t(y) ds_y = \frac{1}{2} g(x) + \int_{\Gamma} \frac{\partial}{\partial n_y} U^*(x, y) g(y) ds_y = f(x)$$

Ansatz (piecewise constant)

$$t(x) \sim t_h(x) = \sum_{k=1}^N t_k \psi_k(x), \quad \psi_k(x) = \begin{cases} 1 & \text{for } x \in \tau_k \\ 0 & \text{elsewhere} \end{cases}$$

Collocation

$$\sum_{k=1}^N t_k \frac{\Delta_\ell}{4\pi} \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y = \Delta_\ell f(x_\ell^*) \quad \text{for } \ell = 1, \dots, N$$

Collocation

$$\sum_{k=1}^N t_k \frac{\Delta_\ell}{4\pi} \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y = \Delta_\ell f(x_\ell^*) \quad \text{for } \ell = 1, \dots, N$$

Galerkin

$$\sum_{k=1}^N t_k \frac{1}{4\pi} \int_{\tau_\ell} \int_{\tau_k} \frac{1}{|x - y|} ds_y ds_x = \int_{\tau_\ell} f(x) ds_x \quad \text{for } \ell = 1, \dots, N$$

Collocation

$$\sum_{k=1}^N t_k \frac{\Delta_\ell}{4\pi} \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y = \Delta_\ell f(x_\ell^*) \quad \text{for } \ell = 1, \dots, N$$

Galerkin

$$\sum_{k=1}^N t_k \frac{1}{4\pi} \int_{\tau_\ell} \int_{\tau_k} \frac{1}{|x - y|} ds_y ds_x = \int_{\tau_\ell} f(x) ds_x \quad \text{for } \ell = 1, \dots, N$$

Stiffness matrices

$$V_h^C[\ell, k] = \frac{1}{4\pi} \Delta_\ell \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y, \quad V_h^G[\ell, k] = \frac{1}{4\pi} \int_{\tau_\ell} \int_{\tau_k} \frac{1}{|x - y|} ds_y ds_x$$

Collocation

$$\sum_{k=1}^N t_k \frac{\Delta_\ell}{4\pi} \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y = \Delta_\ell f(x_\ell^*) \quad \text{for } \ell = 1, \dots, N$$

Galerkin

$$\sum_{k=1}^N t_k \frac{1}{4\pi} \int_{\tau_\ell} \int_{\tau_k} \frac{1}{|x - y|} ds_y ds_x = \int_{\tau_\ell} f(x) ds_x \quad \text{for } \ell = 1, \dots, N$$

Stiffness matrices

$$V_h^C[\ell, k] = \frac{1}{4\pi} \Delta_\ell \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y, \quad V_h^G[\ell, k] = \frac{1}{4\pi} \int_{\tau_\ell} \int_{\tau_k} \frac{1}{|x - y|} ds_y ds_x$$

- dense: N^2 non zero elements

Collocation

$$\sum_{k=1}^N t_k \frac{\Delta_\ell}{4\pi} \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y = \Delta_\ell f(x_\ell^*) \quad \text{for } \ell = 1, \dots, N$$

Galerkin

$$\sum_{k=1}^N t_k \frac{1}{4\pi} \int_{\tau_\ell} \int_{\tau_k} \frac{1}{|x - y|} ds_y ds_x = \int_{\tau_\ell} f(x) ds_x \quad \text{for } \ell = 1, \dots, N$$

Stiffness matrices

$$V_h^C[l, k] = \frac{1}{4\pi} \Delta_\ell \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y, \quad V_h^G[l, k] = \frac{1}{4\pi} \int_{\tau_\ell} \int_{\tau_k} \frac{1}{|x - y|} ds_y ds_x$$

- ▶ dense: N^2 non zero elements \rightarrow Fast Boundary Element Methods
 - ▶ **Fast Multipole Methods** [Greengard, Rokhlin '87, ...]
 - ▶ Panel Clustering [Hackbusch, Nowak '89, ...]
 - ▶ **Adaptive Cross Approximation** [Bebendorf, Rjasanow '03, ...]
 - ▶ Hierarchical Matrices [Hackbusch '99, ...]
 - ▶ Wavelets [Dahmen, Prössdorf, Schneider '93, ...]

Collocation

$$\sum_{k=1}^N t_k \frac{\Delta_\ell}{4\pi} \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y = \Delta_\ell f(x_\ell^*) \quad \text{for } \ell = 1, \dots, N$$

Galerkin

$$\sum_{k=1}^N t_k \frac{1}{4\pi} \int_{\tau_\ell} \int_{\tau_k} \frac{1}{|x - y|} ds_y ds_x = \int_{\tau_\ell} f(x) ds_x \quad \text{for } \ell = 1, \dots, N$$

Stiffness matrices

$$V_h^C[l, k] = \frac{1}{4\pi} \Delta_\ell \int_{\tau_k} \frac{1}{|x_\ell^* - y|} ds_y, \quad V_h^G[l, k] = \frac{1}{4\pi} \int_{\tau_\ell} \int_{\tau_k} \frac{1}{|x - y|} ds_y ds_x$$

- ▶ dense: N^2 non zero elements \rightarrow Fast Boundary Element Methods
 - ▶ **Fast Multipole Methods** [Greengard, Rokhlin '87, ...]
 - ▶ Panel Clustering [Hackbusch, Nowak '89, ...]
 - ▶ **Adaptive Cross Approximation** [Bebendorf, Rjasanow '03, ...]
 - ▶ Hierarchical Matrices [Hackbusch '99, ...]
 - ▶ Wavelets [Dahmen, Prössdorf, Schneider '93, ...]
- ▶ singular surface integrals

Fast boundary element methods: kernel approximation

$$k(x, y) = \frac{1}{4\pi} \frac{1}{|x - y|} \sim \sum_{k=0}^P f_k(x) g_k(y)$$

Fast boundary element methods: kernel approximation

$$k(x, y) = \frac{1}{4\pi} \frac{1}{|x - y|} \sim \sum_{k=0}^P f_k(x) g_k(y)$$

- ▶ Taylor expansion → Panel Clustering

Fast boundary element methods: kernel approximation

$$k(x, y) = \frac{1}{4\pi} \frac{1}{|x - y|} \sim \sum_{k=0}^p f_k(x) g_k(y)$$

- ▶ Taylor expansion → Panel Clustering
- ▶ Spherical harmonics → Fast Multipole

$$k(x, y) = \frac{1}{4\pi} \frac{1}{|x - y|} \approx \frac{1}{4\pi} \sum_{n=0}^p \sum_{m=-n}^n |x|^n Y_n^{-m}(\hat{x}) \frac{Y_n^m(\hat{y})}{|y|^{n+1}}, \quad \frac{|x|}{|y|} < \frac{1}{d}$$

Fast boundary element methods: kernel approximation

$$k(x, y) = \frac{1}{4\pi} \frac{1}{|x - y|} \sim \sum_{k=0}^p f_k(x) g_k(y)$$

- ▶ Taylor expansion → Panel Clustering
- ▶ Spherical harmonics → Fast Multipole

$$k(x, y) = \frac{1}{4\pi} \frac{1}{|x - y|} \approx \frac{1}{4\pi} \sum_{n=0}^p \sum_{m=-n}^n |x|^n Y_n^{-m}(\hat{x}) \frac{Y_n^m(\hat{y})}{|y|^{n+1}}, \quad \frac{|x|}{|y|} < \frac{1}{d}$$

- ▶ Adaptive Cross Approximation

$$s_0(x, y) = 0, \quad r_0(x, y) = k(x, y),$$

Fast boundary element methods: kernel approximation

$$k(x, y) = \frac{1}{4\pi} \frac{1}{|x - y|} \sim \sum_{k=0}^p f_k(x) g_k(y)$$

- ▶ Taylor expansion → Panel Clustering
- ▶ Spherical harmonics → Fast Multipole

$$k(x, y) = \frac{1}{4\pi} \frac{1}{|x - y|} \approx \frac{1}{4\pi} \sum_{n=0}^p \sum_{m=-n}^n |x|^n Y_n^{-m}(\hat{x}) \frac{Y_n^m(\hat{y})}{|y|^{n+1}}, \quad \frac{|x|}{|y|} < \frac{1}{d}$$

- ▶ Adaptive Cross Approximation

$$\begin{aligned} s_0(x, y) &= 0, & r_0(x, y) &= k(x, y), \\ s_k(x, y) &= s_{k-1}(x, y) + \frac{r_{k-1}(x, y_k) r_{k-1}(x_k, y)}{r_{k-1}(x_k, y_k)} \\ r_k(x, y) &= r_{k-1}(x, y) - \frac{r_{k-1}(x, y_k) r_{k-1}(x_k, y)}{r_{k-1}(x_k, y_k)} \end{aligned}$$

Galerkin discretization

$$V_h[j, i] = \int_{\tau_j} \int_{\tau_i} k(x, y) ds_y ds_x$$

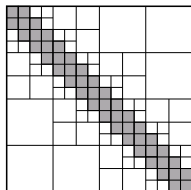
Galerkin discretization

$$V_h[j, i] = \int_{\tau_j} \int_{\tau_i} k(x, y) ds_y ds_x \approx \sum_{k=0}^p \int_{\tau_j} f_k(x) ds_x \int_{\tau_i} g_k(y) ds_y$$

Galerkin discretization

$$V_h[j, i] = \int_{\tau_j} \int_{\tau_i} k(x, y) ds_y ds_x \approx \sum_{k=0}^p \int_{\tau_j} f_k(x) ds_x \int_{\tau_i} g_k(y) ds_y$$

- ▶ Low rank approximation
- ▶ hierarchical clustering
- ▶ admissibility condition

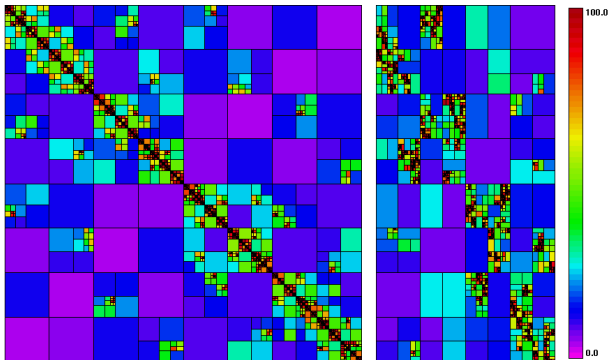


$$\text{dist}(\omega_i^k, \omega_j^k) \geq \eta \max \{ \text{diam } \omega_i^k, \text{diam } \omega_j^k \}$$

- ▶ complexity

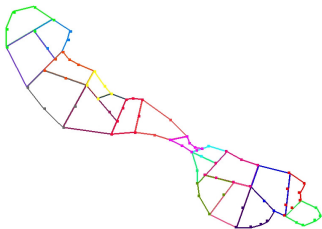
$$\mathcal{O}(N \log^2 N)$$

Adaptive Cross Approximation: Single and double layer potential



non-overlapping domain decomposition

$$-\operatorname{div}[\alpha(x)\nabla u(x)] = 0 \quad \text{for } x \in \Omega, \quad \bar{\Omega} = \bigcup_{i=1}^p \bar{\Omega}_i$$



non-overlapping domain decomposition

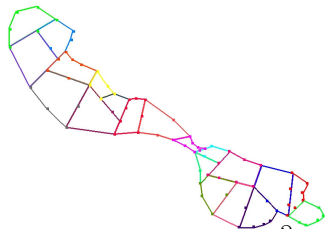
$$-\operatorname{div}[\alpha(x)\nabla u(x)] = 0 \quad \text{for } x \in \Omega, \quad \bar{\Omega} = \bigcup_{i=1}^p \bar{\Omega}_i$$

local boundary value problems

$$\begin{aligned} -\alpha_i \Delta u_i(x) &= 0 & \text{for } x \in \Omega_i, \\ u_i(x) &= g(x) & \text{for } x \in \Gamma_i \cap \Gamma \end{aligned}$$

transmission or coupling boundary conditions

$$u_i(x) = u_j(x), \quad \alpha_i \frac{\partial}{\partial n_i} u_i(x) + \alpha_j \frac{\partial}{\partial n_j} u_j(x) = 0 \quad \text{for } x \in \Gamma_i \cap \Gamma_j$$



non-overlapping domain decomposition

$$-\operatorname{div}[\alpha(x)\nabla u(x)] = 0 \quad \text{for } x \in \Omega, \quad \bar{\Omega} = \bigcup_{i=1}^p \bar{\Omega}_i$$

local boundary value problems

$$\begin{aligned} -\alpha_i \Delta u_i(x) &= 0 & \text{for } x \in \Omega_i, \\ u_i(x) &= g(x) & \text{for } x \in \Gamma_i \cap \Gamma \end{aligned}$$

transmission or coupling boundary conditions

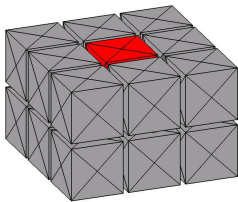
$$u_i(x) = u_j(x), \quad \alpha_i \frac{\partial}{\partial n_i} u_i(x) + \alpha_j \frac{\partial}{\partial n_j} u_j(x) = 0 \quad \text{for } x \in \Gamma_i \cap \Gamma_j$$

local Dirichlet to Neumann map

$$t_i(x) = \frac{\partial}{\partial n_i} u_i(x) = (S_i u_i)(x) \quad \text{for } x \in \Gamma_i$$

- ▶ Local FE/BE approximations of Steklov–Poincaré operators
- ▶ Tearing and interconnecting iterative solution procedures (FETI/BETI)

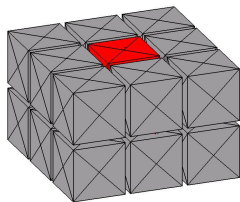
Example: Linear elasticity (steel/concrete)



18 subdomains

L	BETI		all-floating	
	t_2	lt.	t_2	lt.
0	31	19(21(10))	39	22(17(10))
1	217	28(33(14))	170	24(27(14))
2	2129	35(44(14))	1437	27(33(14))
3	14149	42(51(14))	9005	32(36(14))
4	116404	47(54(14))	77111	38(38(15))

Example: Linear elasticity (steel/concrete)



18 subdomains

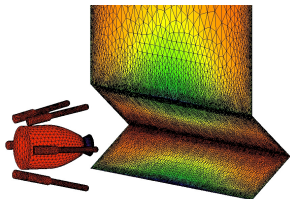
L	BETI		all-floating	
	t_2	lt.	t_2	lt.
0	31	19(21(10))	39	22(17(10))
1	217	28(33(14))	170	24(27(14))
2	2129	35(44(14))	1437	27(33(14))
3	14149	42(51(14))	9005	32(36(14))
4	116404	47(54(14))	77111	38(38(15))

L	N_i	Dirichlet DD		BETI		all-floating	
		t_2	lt.	t_2	lt.	t_2	lt.
0	24	7	53(10)	7	78	8	65
1	96	25	110(14)	19	100	19	82
2	384	181	130(14)	112	114	115	85
3	1536	986	148(14)	562	129	476	95
4	6144	6902	154(14)	4352	153	3119	105
5	24576	59264	166(16)	31645	172	23008	120

24576 boundary elements per subdomain \approx 14 million tetrahedrons

Iterative solution of linear systems

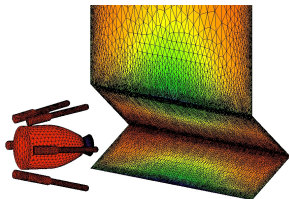
$$\begin{pmatrix} V_h & -K_h \\ K_h^\top & D_h \end{pmatrix} \begin{pmatrix} \underline{t} \\ \underline{u} \end{pmatrix} = \begin{pmatrix} \underline{f}_1 \\ \underline{f}_2 \end{pmatrix}$$



- ▶ scaling
- ▶ preconditioners for V_h , $S_h = D_h + K_h^\top V_h^{-1} K_h$

Iterative solution of linear systems

$$\begin{pmatrix} V_h & -K_h \\ K_h^\top & D_h \end{pmatrix} \begin{pmatrix} \underline{t} \\ \underline{u} \end{pmatrix} = \begin{pmatrix} \underline{f}_1 \\ \underline{f}_2 \end{pmatrix}$$



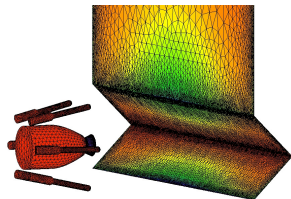
- ▶ scaling
- ▶ preconditioners for V_h , $S_h = D_h + K_h^\top V_h^{-1} K_h$

Preconditioners

- ▶ operators of opposite orders
- ▶ geometric and algebraic multilevel techniques

Iterative solution of linear systems

$$\begin{pmatrix} V_h & -K_h \\ K_h^\top & D_h \end{pmatrix} \begin{pmatrix} \underline{t} \\ \underline{u} \end{pmatrix} = \begin{pmatrix} \underline{f}_1 \\ \underline{f}_2 \end{pmatrix}$$



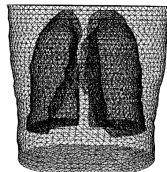
- ▶ scaling
- ▶ preconditioners for V_h , $S_h = D_h + K_h^\top V_h^{-1} K_h$

Preconditioners

- ▶ operators of opposite orders
- ▶ geometric and algebraic multilevel techniques

Acoustic and electromagnetic scattering problems

- ▶ Helmholtz, Maxwell
- ▶ combined boundary integral equations to avoid spurious modes
- ▶ domain decomposition methods, preconditioners



Optimization and inverse problems, eigenvalue problems, multiphysics, fluid structure interaction, FEM/BEM coupling, ...

Some Activities

- ▶ Söllerhaus Workshops on Fast Boundary Element Methods in Industrial Applications. Kleinwalsertal, Austria.
September 30–October 3, 2010
September 29–October 2, 2011
- ▶ IABEM Symposium, Brescia, Italy, September 5–8, 2011
- ▶ Annual GAMM Meeting, Graz, Austria, April 18–21, 2011

Some References

- ▶ O. Steinbach: Numerical Approximation Methods for Elliptic Boundary Value Problems. Finite and Boundary Elements. Springer, New York, 2008.
- ▶ S. Rjasanow, O. Steinbach: The Fast Solution of Boundary Integral Equations. Mathematical and Analytical Techniques with Applications to Engineering. Springer, New York, 2007.
- ▶ M. Schanz, O. Steinbach (eds.): Boundary Element Analysis. Mathematical Aspects and Applications. Lecture Notes in Applied and Computational Mechanics, vol. 29, Springer, Heidelberg, 2007.
- ▶ U. Langer, M. Schanz, O. Steinbach, W. L. Wendland (eds.): Fast Boundary Element Methods in Engineering and Industrial Applications. Lecture Notes in Applied and Computational Mechanics, Springer, Heidelberg, 2011.

<http://www.numerik.math.tu-graz.ac.at>